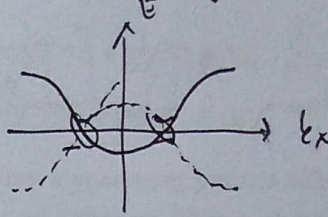
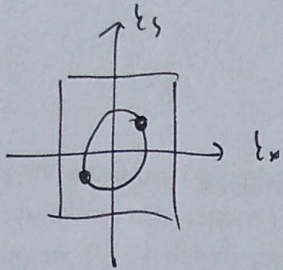


Andreev reflection

A 1

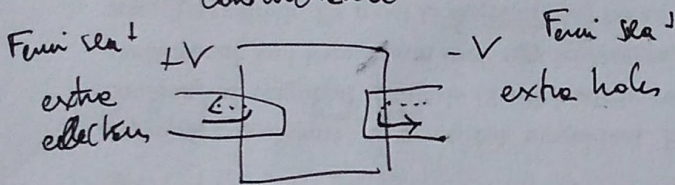
1) Formation of a Cooper pair = incoming electron → reflected hole



2) Hole left by taking e^- @ $+k, \uparrow$, right-mover:

Andreev reflection keeps k constant! \downarrow , right-mover

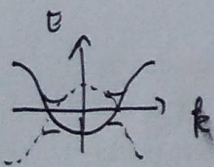
3) Inverse Andreev: hole → electron. Takes place e.g. during constrictance



3) Andreev reflection is coherent. Amplitude obtained by BTK

$$k_F : \begin{cases} r_A = e^{-i\varphi} e^{-i \arccos \frac{E}{\Delta}} \\ -k_F : \tilde{r}_A = e^{+i\varphi} e^{-i \arccos \frac{E}{\Delta}} \end{cases}$$

s-wave: $\hat{H} = \begin{pmatrix} \tilde{c}_{k\uparrow} & \tilde{c}_{-k\downarrow} \\ \Delta^\dagger(x) & 2t\cos k + \mu \\ \Delta^\dagger(x) & 2t\cos k + \mu \end{pmatrix} \begin{pmatrix} \tilde{c}_{k\uparrow} \\ \tilde{c}_{-k\downarrow} \end{pmatrix}$



(above gap: HW)

- linear H = $\begin{bmatrix} -iv_F \partial_x & \Delta \\ \Delta^\dagger & iv_F \partial_x \end{bmatrix}$

- EFA

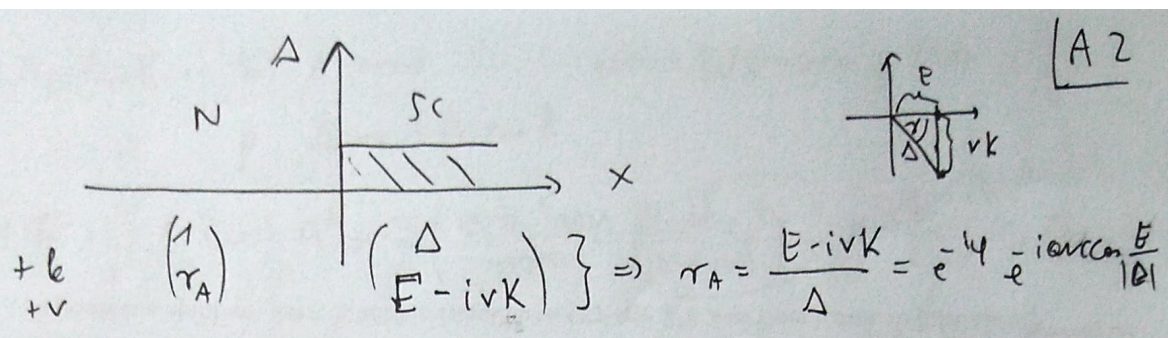
Inside gap: $e^{i\epsilon_F x} e^{-Kx}$

$$H = \begin{bmatrix} iv_k & \Delta \\ \Delta^\dagger & -iv_k \end{bmatrix} = iv_k \sigma^z + \text{Re} \Delta \sigma^x - \text{Im} \Delta \sigma^y \quad \Delta = \Delta e^{i\varphi}$$

$$H^2 = \sqrt{4\Delta^2 - v^2 k^2} \quad v_k = \pm \sqrt{4\Delta^2 - E^2}$$

need only + energy: $\begin{pmatrix} \Delta \\ E - iv_k \end{pmatrix} = \begin{pmatrix} u \\ v \end{pmatrix}$ penetration depth diverges @ $\Delta = E$

$|u|^2 = |v|^2$ electron-hole weight equal



Inverse Ansatz:

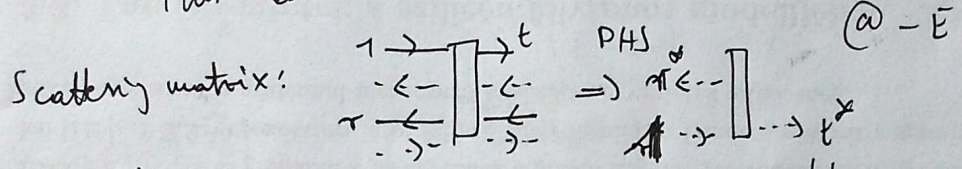
$$\begin{pmatrix} 1 \\ r_A \end{pmatrix} \begin{pmatrix} \Delta \\ E - ivk \end{pmatrix} \Rightarrow r_A = \frac{E - ivk}{\Delta} = e^{-i\varphi} e^{-i \arccos \frac{E}{|\Delta|}}$$

$$\begin{pmatrix} \tilde{r}_A \\ 1 \end{pmatrix} \begin{pmatrix} \Delta \\ E + ivk \end{pmatrix} \Rightarrow \tilde{r}_A = \begin{pmatrix} \Delta \\ E + ivk \end{pmatrix} = e^{+i\varphi} e^{-i \arccos \frac{E}{|\Delta|}}$$

other valley

4) Conductance across a tunnel junction weaker in SC

This can be used to measure gap.



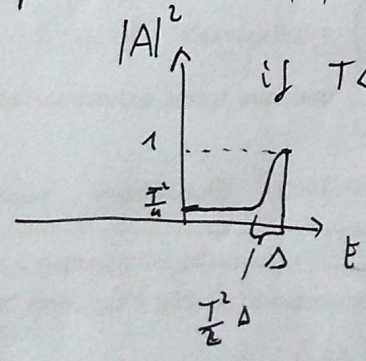
Calculate Ansatz:

Naive: $r_A \sim |t|^2$

Correction due to multiple scattering:

$$A = t^{1*} r_A [1 - r_A \tilde{r}_A r_A^{1*} r_A] t$$

single mode: $|A| \approx \left| r_A \frac{t^{1*} t}{1 - e^{-2i \arccos \frac{E}{\Delta}} |r'|^2} \right| = \frac{T}{1 - e^{-2i \arccos \frac{E}{\Delta}} R}$



if $T \ll 1$: $A \approx e^{-2i \arccos \frac{E}{\Delta}} R \gg T$

$E = 0$: $|A| = \frac{T}{1+R} = \frac{T}{2-T}$

$E = |\Delta|$: $|A| = \frac{T}{1-R} = 1$

$1 - e^{-2i \arccos \frac{E}{\Delta}} R \gg T$ except $\frac{v k}{\Delta} < T$

front

$$\left. \begin{aligned} v^2 k^2 < T^2 \Delta^2 \\ \Delta^2 - E^2 \end{aligned} \right\} \Delta - E < \frac{T^2}{2} \Delta$$

$$\beta + E (\Delta - E) \approx 2\Delta (\Delta - E)$$

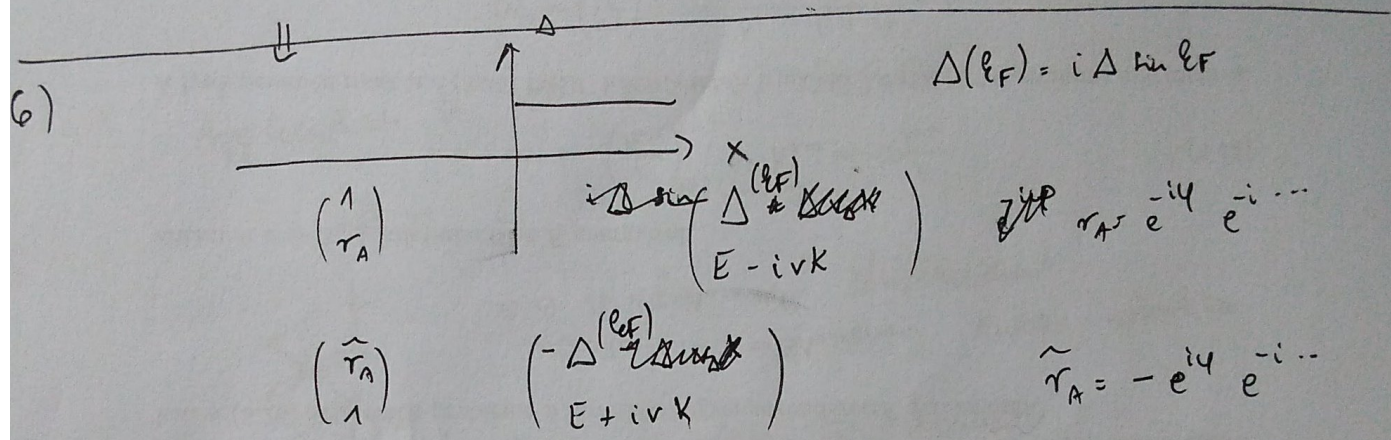
5) Application to p-wave SC: use ~~BdG~~ via ρ BldG A

$\rightarrow \rho \Delta \rightarrow i \Delta \sin k$

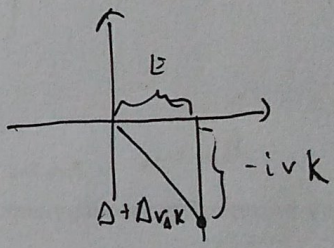
$\hat{H} = \sum_k [t + 2 \cos k c_k^\dagger c_k - \mu c_k^\dagger c_k] + \Delta \sum_k e^{ik} c_k^\dagger c_k^\dagger + \Delta \sum_k e^{-ik} c_k c_k$

\Downarrow only $\sin k$ only $\sin k$ remains

$\mathcal{H}(k) = \begin{bmatrix} -t \cos k - \mu/2 & i \Delta \sin k \\ -i \Delta \sin k & t \cos k + \mu/2 \end{bmatrix}$ $E = \sqrt{\Delta^2 (1 + v_\Delta k)^2 - v^2 k^2}$



$|A| = \frac{T}{1 + e^{2i\dots} \cdot R}$



@ $E=0$: ~~$\sqrt{2} |A|^2 v_\Delta^2 k$~~

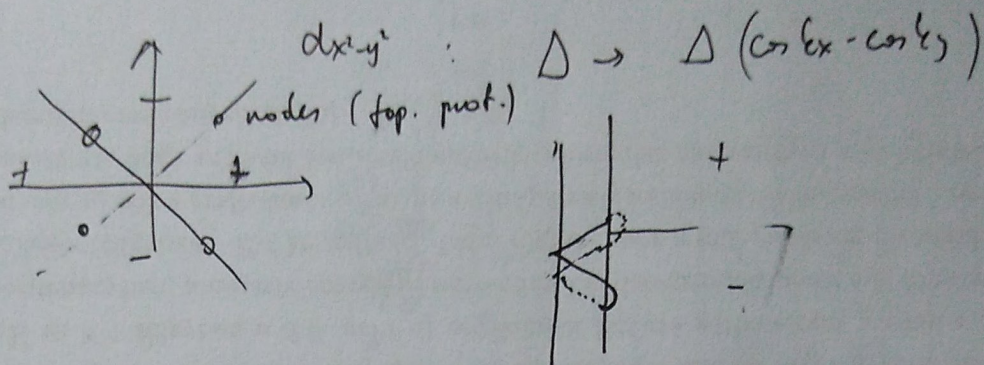
Narrow resonance, $\frac{E^2}{v^2 k^2} < \Delta^2$

Narrow resonance around $E=0$: $|A|=1$ ZBCP

@ $E=\Delta$: $k=0$

$|A| = \frac{T}{1+R}$

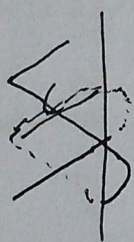
7) Application to 2D d-wave (cuprates): ZBCP



depending on angle of surface

8) Application to 2D p+ip (Sr_2CuO_4 ?)

$$\Delta \rightarrow \Delta \sin k_x + i \sin k_y$$



$\varphi \rightarrow \varphi \cdot k \cdot 1 \rightarrow$ displacement during Andreev reflection

same direction displacement

@ were Andreev

edge state!

Chern # \Rightarrow Propagating Majorana modes